product is thermally unstable and isomerizes under the reaction conditions to produce what may well be the more stable *cis* product. A second route to the observed product involves collapse of the original *trans*octahedral structure to a trigonal-bipyramidal intermediate. This intermediate is considered to have the fluoro ligand in an apical position and one of the ethylenediamine ligands in the trigonal plane. The remaining apical and trigonal positions are occupied by the second ethylenediamine molecule.

After the collapse of the octahedron the entering anion attacks in the trigonal plane of the intermediate to produce a *cis* product. At the present time it is not possible to prove or disprove either of the proposed routes to the final product. However, we are currently investigating various synthetic approaches to complexes of the type $trans-[Cr(en)_2FX]^n$ ⁺ which should shed some light on the nature of the intermediates involved in the synthesis of the *cis* complexes.

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The Crystal and Molecular Structure of the High- Spin Square-Planar Complex Triphenylmethylarsonium Bis (toluene- 3,4-di thiola to)cobal ta te-0.5 - **E than01**

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The crystal and molecular structure of the high-spin complex triphenylmethylarsonium bis(toluene-3,4-dithiolato)cobaltate-0.5-ethanol, $[(C_6H_5)_8(CH_3)As][Co(tdt)_2]~0.5C_2H_6OH$, has been determined from three-dimensional single-crystal X-ray data collected by counter methods. The structure has been refined by least-squares procedures to a conventional *R* factor of 0.063 for 1108 nonzero reflections. The complex crystallizes in space group \overline{PI} of the triclinic system with two molecules in a unit cell of dimensions $a = 18.61 \pm 0.01$, $b = 10.848 \pm 0.006$, $c = 10.206 \pm 0.005$ Å, $\alpha = 112.45 \pm 0.02^{\circ}$, $\beta = 71.55 \pm 0.02^{\circ}$ 0.02°, $\gamma = 115.90 \pm 0.02$ °, and $V = 1685.4 \text{ Å}^3$. An observed density of 1.359 \pm 0.010 g/cm³ is in excellent agreement with the calculated value of 1.355 g/cm³ for $Z = 2$. The Co(tdt)₂⁻ anions are crystallographically required to have centers of symmetry since the two cobalts per unit cell occupy the $(0, 0, 0)$ and $(1/2, 0, 1/2)$ special positions of the space group. The S₄ coordination about the cobalts is very nearly square planar with average intra- and interligand *S*-S distances of 3.098 and 3.028 Å, respectively, and estimated standard deviations for these values of 0.007 Å. The planar Co(tdt)₂⁻ anions have an ordered *trans* configuration of the tdt ligand methyl groups and thus possess C_{2b} symmetry. The $(C_eH_8)_8(CH_3)As^+$ cation has its expected tetrahedral geometry about arsenic. Ethanol solvent molecules were trapped during crystallization of the complex, presumably because of the asymmetric nature of the cation. These solvent molecules are disordered about the $(0, \frac{1}{2})$, 0) center of symmetry with only one molecule per unit cell. Thus, the crystal structure of the complex consists of the packing of essentially planar anions, bulky cations, and trapped solvent molecules. The anions are well separated and exhibit none of the magnetic interactions which characterize the structures of the other related mononegative bis complexes. A ground-state electronic structural assignment for the spin-triplet $Co(tdt)_2$ ⁻ complex is proposed on the basis of molecular orbitals calculated for a closely related nickel complex using the observed structural parameters.

Introduction

The bidentate sulfur-donor ligand systems of types I and I1 have been found to be excellent stabilizers of the square-planar geometry for a wide variety of metals and ground-state configurations.² Although we show I and **I1** in their classical dianion formulation, it is

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clear from electronic and molecular structural studies that these ligand systems coordinate to form unsaturated five-membered metal-chelate rings in which π bonding plays an important stabilizing role. Complexes of the four ligand systems maleonitriledithiolate (mnt, I with $Q = CN$), *cis*-stilbenedithiolate (S₂C₂- $(C_6H_5)_2$, $cis-1,2-bis(trifluorometry$) ethylenedithiolate $(S_2C_2(CF_3)_2)$, and toluene-3,4-dithiolate (tdt, II with $X = Y = H$, $Z = CH₃$) have received particular attention in these studies.

The considerable interest and activity which de- **(2) Four of the more recent papers dealing with the synthesis, characterization, and electronic structural properties of the his square-planar complexes and their oxidation-reduction reactions are: (a) G.** N. **Schrauzer** and V. P. Mayweg, *J. Am. Chem. Soc.*, 87, 3585 (1965); (b) R. Williams, E.
Billig, J. H. Waters, and H. B. Gray, *ibid.*, 88, 43 (1966); (c) M. J. Baker-**Hawkes, E. Billig, and H. B. Gray,** *ibid.,* **88, 4870** (1966). **(d)** A. **Davison, D.** V. **Howe, and E. T. Shawl,** *Inorg. Chem.,* **6, 458 (1907). References to earlier studies in this field are given therein.**

veloped in the bis complexes of the bidentate sulfurdonor ligand systems have prompted several detailed structural investigations on them.3-8 Most of the structure determinations have been performed on members of the bis(maleonitriledithio1ate) series including the dianion of Ni ,³ the dianion of Co ,⁴ the monoanion of $Cu,^5$ and the monoanion of Ni.⁷ The coordination geometry surrounding the metals in these complexes is essentially square-planar and the point group symmetry of all of the anions is very nearly D_{2h} . Despite the differences that are found to exist in the packing of the molecules, the mnt series appears quite remarkable since no significant variations are observed in the geometry and dimensions of the anions in all four structures. The molecular structure of the neutral complex $Ni(S_2C_2(C_6H_5)_2)$ has also been reported⁸ and its coordination geometry is essentially identical with that found for the bis(mnt) anions, Recent studies, however, have demonstrated that certain members of the bis complexes of Co and Fe exist as five-coordinate dimers in which the metal is surrounded by a square pyramid of sulfur atoms.^{6,9} In this paper and one to follow, 10 we present the structure determinations of the two closely related monoanionic complexes bis(to1 uene-3,4-dithiolato)cobaltate¹¹ and bis(3,4,5,6-tetra**chlorobenzene-1,2-dithiolato)cobaltate.2c** These detailed structural studies play an important role in evaluating descriptions of the electronic structures of these bis complexes and in elucidating the factors which influence their tendency toward dimerization.

The monoanion $Co(tdt)₂$ is probably the most extraordinary of the bis complexes of the bidentate sulfur-donor ligand systems. The complex was first synthesized in 196311 and found to have a stable spintriplet ground state in solid salts and in solutions utilizing a wide variety of solvents. Therefore, the complex $Co(tdt)₂-$ is the first apparent example of a square-planar transition metal complex possessing a high-spin ground-state configuration. A detailed single-crystal X-ray investigation was therefore undertaken in order to confirm the apparent square-planar coordination for the $Co(tdt)₂$ anion and to establish structural parameters for electronic structural studies of this unusual system. In addition, we believed it important to establish the packing of the $Co(tdt)₂-$ anions in the crystal since structural studies of the related complexes $[(C_6H_5)_8(CH_3)P][Ni(mnt)_2]^7$ and $[(n-C_4H_9)_4N]$ - $[Cu(mnt)₂]$ ⁵ revealed the planar anionic units to lie in relatively close proximity to one another. Spin-spin interactions which may thus occur in the crystalline

(7) C. J. Fritchie, *Acta Cryst., 20,* 107 (1966).

state have been proposed as an important factor in the magnetic properties of the paramagnetic $Ni(mnt)_2$ complex.12

The structure determination of the $(C_6H_5)_3(CH_3)As^+$ salt of the $Co(tdt)₂$ complex anion proved considerably more difficult than was originally anticipated. Although a set of integrated precession intensity data was initially collected in 1964, the problem defied solution for reasons which became apparent when it was reinvestigated and solved in 1966. The space group and unit cell determinations along with certain preliminary considerations are taken from the work at Brookhaven National Laboratory. A refinement of the cell constants, the collection of a new set of intensity data, and the solution of the structure were carried out at Columbia University.

Unit Cell and Space **Group** Determination

A generous supply of dark blue crystals of $[(C_6H_5)_{3-}]$ $(CH_3)As][Co(tdt)_2]$ were kindly given to us by Dr. Raymond Williams.2b The crystals were examined optically and by Weissenberg and precession photography and found to belong to the triclinic system. A convenient primitive cell was chosen with the dimensions $a = 18.60, b = 11.01, c = 10.32 \text{ Å}, \alpha = 114.2^{\circ},$ $\beta = 71.0^{\circ}, \gamma = 116.7^{\circ}, \text{ and } V = 1695 \text{ Å}^3.$ A Delaunay reduction performed on this cell failed to reveal the presence of any symmetry. The reduced cell parameters which were not used in any subsequent calculations are $a = 18.09$, $b = 16.82$, $c = 10.32$ Å, $\alpha = 95.3^{\circ}, \beta = 103.6^{\circ}, \text{ and } \gamma = 141.4^{\circ}.$ No piezoelectric effect was looked for in the unit cell determination. An experimental density of 1.36 ± 0.03 g/cm3 determined by pycnometry was in reasonable agreement with a value of 1.35 g/cm^3 calculated for two $[(C_6H_5)_3(CH_3)As][Co(tdt)_2]$ molecules using the preliminary cell constants given above. On the basis of these results, we assumed the centrosymmetric space group \overline{PI} for the complex. This assumption implies that no symmetry conditions need be imposed on the molecule and all of its atoms may occupy general positions of the space group. Although this was found not to be the case, the satisfactory agreement ultimately obtained between observed and calculated structure factors strongly suggests that \overline{PI} is the correct choice.

Collection and Reduction of the Intensity Data

A complete set of intensity data mas collected from a single crystal at room temperature using a Picker diffractometer fitted with a General Electric quartercircle cradle. The crystal was mounted in the $\overline{a^*}$ direction on a eucentric goniometer head and aligned ac. curately on a Supper precession camera. After transferring the crystal to the goniostat, the direction of a^* was adjusted to coincide with the polar axis of the diffractometer by centering the $\bar{k}00$ reflections in a diffracting position at $\chi = 90^{\circ}$ for all values of the polar or spindle angle ϕ . Approximately 40 reflections.

⁽³⁾ R. Eisenberg, *3.* A. Ibers, R. J. H. Clark, and H. B. Gray, *J. Am. Chem. SOL, 86,* 113 (1964); R. Eisenberg and **5. A.** Ibers, *Inoug. Chem.,* **4,** *605* (1965).

⁽⁴⁾ J. D. Forrester, A. Zalkin, and D. H. Templeton, *ibid.,* **3,** 1500 (1964).

⁽⁵⁾ J. D. Forrester. A. Zalkin, and D. H. Templeton, *ibid.,* **S,** 1507 (1964).

⁽⁶⁾ J. H. Enemai-k and W. N. Lipscomb, *ibid.,* **4,** 1729 **(1965).**

⁽⁸⁾ D. Sartain and M. R. Truter, *Chem. Commun.*, 382 (1966).
(9) (a) A. L. Balch and R. H. Holm, *ibid.*, 552 (1966); (b) G. N. Schrauzer, V. P. Mayweg, H. W. Finck, and W. Heinrich, *J. Am. Chem. Soc.*, 88, 4604 (1966); *(c)* W. *C.* Hamilton and I. Bernal, *Inorg. Chem.,* 6,2003 (1967).

⁽¹⁰⁾ M. J. Baker-Hawkes, Z. Dori, R. Eisenberg, and H. B. Gray, J. *Am. Chem. Sor.,* in press.

⁽¹¹⁾ H. B. Gray and E. Billig, *ibid.,* **85.** 2019 (1963).

of both general and special classes were then centered in the counter aperture by varying 2θ , χ , and ϕ in conjunction with the left-right and top-bottom balancing features of the variable receiving aperture. After this had been accomplished, the orientation and unit cell parameters were refined by a least-squares procedure using W. C. Hamilton's mode **1** diffractometer setting program.¹³ The refinement yielded the following values of the cell constants and their estimated standard deviations: $a = 18.61 \pm 0.01$, $b = 10.848 \pm 0.006$, $c = 10.206 \pm 0.005$ Å, $\alpha =$ $112.45 \pm 0.02^{\circ}, \ \beta = 71.55 \pm 0.02^{\circ}, \ \gamma = 115.90 \pm 0.02^{\circ}$ 0.02°, and $V = 1685.4 \text{ Å}^3$. The mosaic spread of the crystal was then determined from open source-narrow counter ω scans¹⁴ and found to be relatively isotropic with an average value of **0.30".**

The intensities of **1142** independent reflections with an angular range $\theta_{Cu} \leq 35^\circ$ were measured by the stationary crystal-stationary counter technique using nickel-filtered Cu $K\alpha$ radiation. Only a small number of reflections beyond $\theta = 35^{\circ}$ appeared to have significant intensity. Since these reflections are composed primarily of contributions from the heavy atoms, their intensities were not determined. The peak heights of the measured reflections were counted for 20-sec durations while the background was measured in 10-sec counts both above and below the **20** value for each reflection. The data were collected in two arbitrary segments, each being assigned its own scale factor. In processing the data, we considered a reflection to be unobserved if the peak count did not exceed the combined background counts for the reflection. The usual Lorentz-polarization factors were applied to the observed intensities to yield a set of F_0^2 values, where *F,* is the observed structure factor amplitude. An absorption correction was then carried out for the observed structure amplitudes using the GONO9 diffractometer absorption program. The six crystal faces were first identified by optical goniometry and their dimensions carefully measured. With an absorption coefficient of **73.4** cm-' and approximate crystal dimensions of $0.36 \times 0.19 \times 0.08$ mm, the resultant transmission coefficients were found to range from **0.23** to 0.58 . The F_0 values were then brought to an approximate absolute scale through a modification of Wilson's method.

Solution and Refinement of the Structure

Because of the assumption of space group \overline{PI} with two molecules in the unit cell and the rather asymmetric nature of the $(C_6H_5)_3(CH_3)As^+$ cation, it was presumed that all atoms of the molecule occupy the 2(i) general positions of the space group $\pm [x, y, z]$.¹⁵ This presumption requires that the two $Co(tdt)_2$ anions are symmetry related through a center of inversion. If, in addition, the complex anions are very nearly square planar, one would expect to see a highdensity Patterson peak corresponding to the intermolecular Co-Co vector which is also reinforced by intermolecular S-S vectors and C-C vectors. In addition, the presence of As in the unit cell yields doubly degenerate sets of As-Co and As-S vectors together with one nondegenerate As-As vector.

Numerous models based on a computed three-dimensional Patterson function were formulated in accordance with the assumption of general positions, but none of these models was successfully refined. After a careful reexamination of the Patterson map a completely consistent set of vectors was derived in which the two cobalts per unit cell were located at the centers of symmetry $(0, 0, 0)$ and $(\frac{1}{2}, 0, \frac{1}{2})$ while the two arsenics remained in general positions of the space group. The positions of the four independent sulfur atoms were then quickly determined from the Patterson function. The positional parameters of the S and As atoms along with variable isotropic temperature factors assigned to each of them and the two *CO'S* were refined through several cycles of least squares. From a subsequent difference Fourier based on phases obtained from these heavy atoms, **11** of the **12** anion benzene ring carbons and all **19** of the cation carbons were located in very reasonable positions. No attempt was made at this stage of refinement to find the methyl groups of the tdt ligands because of expected disorder in their positions.

The trial structure was refined using the JGLS least-squares group refinement procedure in which the phenyl rings were constrained to their well-known geometry (D_{6h} symmetry, C-C = 1.392 Å, C-H = 1.08 Å). The function minimized was $\Sigma w (F_o - F_o)^2$ where the weights w of the reflections were set equal to the squares of the reciprocals of the corresponding standard deviations σ . The function $(TLp)^{-1}[CT + CB + (0.03(CT - CB))^2]^{1/2}/2F_0$ was used to estimate the value of σ for each reflection where CT is the total count, **CB** is the total background count, *Lp* is the Lorentz-polarization factor, and *T* is the transmission coefficient for the particular reflection.16 The neutral Co, As, S, and C scattering factors were taken from Ibers' tabulation¹⁷ while the anomalous parts of the Co, As, and S scattering factors were obtained from Templeton's tabulation¹⁸ and included in the calculated structure factors.¹⁹

In the refinement, each phenyl ring was assigned a single variable isotropic thermal parameter and six variable positional parameters which include the three

⁽¹³⁾ In addition to the mode 1 setting program, the main programs for the IBM 7094 used in this work were Hamilton's GONO9 diffractometer **absorption program, the Busing-Levy ORFLS least-squares program, the Zalkin FORDAP Fourier program, the ORFPB function and error program, and the JGLS group refinement least-squares program written by J. A. I. Various other programs from the Brookhaven Crystallographic Library were also used in the investigation. Most of the computation was performed on Columbia's 7094.**

⁽¹⁴⁾ T. C. Furnas, "Single Crystal Orienter Instruction Manual," General Electric Co., Milwaukee, Wis., 1967, Chapter 10, p 138.

^{(15) &}quot;International Tables for X-Ray Crystallography," The Kynoch Press, Vol. I, Birmingham, England, 1965, p 75.

⁽¹⁶⁾ W. R. Busing and H. A. Levy, *J. Chem. Phys.,* **26, 563 (1957). (17)** J. **A. Ibers, "International Tables for X-Ray Crystallography,"**

Vol. **3, The Kynoch Press, Birmingham, England, 1965, Table 3.3.1A.**

⁽¹⁸⁾ D. H. **Templeton, "International Tables for X-Ray Crystallogra phy,"** Val. **3, The Kynoch Press, Birmingham, England, 1965, Table 3.3.2B.**

⁽¹⁹⁾ J. A. Ibers and W. C. Hamilton, *Acta* **Crysf., 17, 781 (1964).**

fractional coordinates of the ring center and three rotation angles describing the ring orientation. The initial refinement of the cation phenyl rings together with the 19 nongroup atoms, each of which had a variable isotropic temperature factor, converged to a coniable isotropic temperature factor, converged to a conventional *R* factor $(R = \Sigma ||F_o| - ||F_o|| / \Sigma |F_o|)$ of to ventional *R* factor $(R = \Sigma ||F_o| - ||F_e|| / \Sigma |F_o|)$ of 0.148 and a weighted *R* factor *R'* $(R' = \langle \Sigma w(F_o F_e$)²/ $\sum w F_o$ ²)^{1/2}) of 0.168 for the 93 positional, scale, and thermal parameters. A difference Fourier based on this refinement showed in the asymmetric unit six peaks with electron density greater than 1.2 e/ \AA ³ or approximately 50% of the height of a carbon atom in this structure. One of these peaks was easily identified as the missing carbon atom in the benzene ring of one of the anions while two others were identified as the ordered methyl groups of the tdt ligands. The three remaining peaks which were located more than 3.5 Å from either the cations or anions were assigned as the atoms of a trapped molecule of ethanol from which the crystals had been grown. These six atoms were included in the next round of isotropic refinements but no attempt was made to distinguish which of the solvent atoms was oxygen. This refinement of 117 positional, scale, and thermal parameters converged to conventional and weighted *R* factors of 0.093 and 0.120 for 1113 nonzero reflections.

A difference Fourier based on this refinement revealed electron density as high as $0.7 \text{ e}/\text{\AA}$ near the arsenic atoms and as high as $0.6 \text{ e}/\text{\AA}^3$ in regions near the anions, thus providing evidence for some anisotropic vibration of the heavy atoms. The isotropic refinement also led to temperature factors for the solvent molecule ranging from 12 to 20 \AA ². This was the first indication that the solvent molecule sites in the unit cell may not be fully occupied. A subsequent redetermination of the density and an elemental analysis later confirmed this notion.²⁰

After the elimination of several indexing and punching errors, a further refinement was carried out in which the Co, As, and S atoms were allowed to vibrate anisotropically, while all other atoms were restricted to isotropic vibration. In addition, fixed contributions from the phenyl ring hydrogens were added to the calculatcd structure factors. This final refinement of 150 positional, thermal, and group parameters converged to the final conventional and weighted *K* factors of 0.063 and 0.077, respectively. With the weighting scheme employed, the estimate of the standard deviation of an observation of unit weight has the value 2.77. The final agreement between F_0 and $|F_c|$ for reflections of high intensity appears good, thus indicating the reasonableness of the 0.03 factor used in the estimation of $\sigma(F^2)$. A difference Fourier based on this refinement showed no peaks higher than 0.45 e/ \AA ³ or about 20% of the height of a carbon atom, and reasonable positions were found for 12 of the remaining 18 hydrogen atoms in the asymmetric unit. The contributions of these additional hydrogens to $|F_c|$ were not included in any additional calculations.

The parameters obtained in this final round of calculations are taken as the final parameters for the structure and are given in Table I. Standard deviations, as obtained from the inverse matrix, are also given. Table I1 lists the fractional coordinates of the cation phenyl ring carbon atoms which can be derived from the group parameters in Table I.

The F_0 and $|F_0|$ values (in electrons \times 10) for the 1108 observed reflections included in the final refinement are tabulated in Table III. The $|F_e|$ values for the unobserved reflections are omitted from this table.

The anisotropic thermal parameters can be interpreted in terms of the real amplitudes of vibration, since they are essentially uncorrelated with the scale factors. It should also be noted that in the partial anisotropic refinement, 35 carbon atoms including the phenyl ring carbons and the solvent atoms were restricted to isotropic vibration. In Table IV, the principal values of the root-mean-square amplitudes of vibration are listed. The orientations of the thermal ellipsoids, which may be derived from the parameters in Table I, are shown in Figure 1.

In order to check the correctness and accuracy of the initial set of integrated precession data, a final structure factor calculation was performed using the parameters of Table I. Only the scale factors for the different levels and zones were adjusted. This calculation yielded an agreement factor of $R = 0.143$ between the 998 structure amplitudes from the precession data and the final calculated structure factors. The precession data, which were collected using Mo K_{α} radiation, were uncorrected for absorption. Thus, the precession data which were of good quality could have led to essentially the same structure.

Description of the Structure

The location of the Co atoms at centers of symmetry removes the constraint that all of the anions in the crystal must be parallel but necessitates the immediate coordination geometry around the cobalts to be planar. The \cos_4 geometry is, in fact, very nearly square

⁽²⁰⁾ In order to estahlish clearly whether or not an ethanol molecule was present in the asymmetric unit, the density was carefully redetermined by pycnometric methods and yielded an experimental value of 1.359 ± 0.010 g/ems, Densities were then calculated for several molecular formulas with the use of the refined cell constants. These calculated values are: (a) $[(C_6H_5)_3(CH_3)As][Co(tdt)_2], \rho_0 = 1.312 g/cm^3;$ (b) $[(C_6H_5)_3(CH_3)As]$ $[Co(tdt)_1] \cdot C_2H_5OH$, $\rho_c = 1.399$ g/cm³; (c) $[(C_6H_5)_3(CH_3)As][Co(tdt)_2]$. $0.5C_2H_6OH$, $\rho_0 = 1.355$ g/cm³. Clearly, the values for (a) and (b) are in significant disagreement with the experimental density and strongly indicate the improbability of either molecular formula. On the other hand, the calculated density for the hemiethanolate *(c)* is in excellent agreement with the experimental value of 1.359 g/cm³. Crystals from the sample used in the X-ray investigation were then submitted for elemental analysis, including a direct oxygen determination, to Midwest Microlab., **Inc.** *Anal.* Calcd for **[(CaHa)s(CHa)As][Co(tdt)n]:** C, **57.56;** H, **4.36;** 0, *0.00;* S, **18.60.** Calcd for **((CsH6)~(CHs)As][Co(tdt)~]** 'C?HsOH: C, **57.22;** H, **4.90;** *0,* **2.18;** S, **17.44.** Calcd for $[(C_6H_8)_8(CH_8)AS][Co(tdt)_2] \cdot 0.5C_2H_8OH: C, 57.38; H, 4.64; O, 1.13; S, 18.00. Found: C, 57.65; H, 4.66; O, 1.23; S, 17.75.$ The analytical data thus appear to favor the hemiethanolate formula, and the direct determination of **1,23'%** oxygen provides positive evidence that trapped ethanol molecules are present in the crystal. The density and analytical results strongly support conclusions based on an examination of the finaldistances of the complex which indicates that only one ethanol molecule can be present in each unit cell. This examination and conclusions based **on** it are discussed in the text. The very satisfactory agreement which **was** ultimately obtained in the refinement is taken as further confirmation of the correctness of the hemiethanolate assignment.

TABLE I FINAL PARAMETERS FOR $[(C_6H_5)_3(CH_3)As][C_0(tdt)_2] \cdot 0.5C_2H_5OH$ Atomic Positional and Isotopic Thermal Parameters

a The numbers in parentheses here and in succeeding tables are estimated standard deviations in the least significant digits. $\frac{1}{2}X(1)$, $X(2)$, and $X(3)$ refer to the atoms of the trapped ethanol molecule which is discussed in the text. The positional and thermal parameters given above are taken from the final round of calculations in which the multiplicity factors of these atoms were held fixed at 1.0 and all three atoms were assigned carbon atom scattering factors. An additional calculation in which $X(1)$, $X(2)$, and $X(3)$ were assigned fixed multiplicity factors of 0.5 and $X(1)$ was assumed to be the ethanol oxygen yielded the following parameters for the three atoms.

c The form of the anisotropic thermal ellipsoid is: $\exp[-(\beta_{11}h^2 + \beta_{21}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)]$. $d x_0, y_0$, and z_0 are the fractional coordinates of the ring centers. The angles ϕ , θ , and ρ (in radians) are the ring orientation angles which bring about alignment (except for translation) of an internal coordinate system within the ring with a fixed external coordinate system which have been defined in previous papers. See, for example, R. Eisenberg and J. A. Ibers, *Inorg. Chem.,* 4,773 (1965), or S. J. La Placa and J. A. Ibers, *Acta Cryst.,* 18, 511 (1965).

planar with average intra- and interligand S-S distances of 3.098 and 3.028 A, respectively, and estimated standard deviations for these values of 0.007 Å . Principal intramolecular distances and angles for the complex are tabulated in Table V. In addition, the two $Co(tdt)₂$ anions are shown with their dimensions in Figure 1. The average Co-S and S-C bond lengths of 2.166 and 1.78 *k* agree reasonably well with the corresponding distances which have been found in the related bis(mnt) structures $(2.161$ and 1.72 Å in Co- $(\text{mnt})_2^2$, 4 2.146 and 1.71 Å in Ni $(\text{mnt})_2$, 7 2.166 and 1.76 Å in $Ni(mnt)₂^{2-3}$ and 2.170 and 1.72 Å in Cu- $(mnt)_2 - 5$.

Both anions are essentially planar and possess C_{2h} symmetry although only C_i symmetry is required. The Co atoms necessarily sit on the best weighted leastsquares planes through the entire anion units. In Table VI, the equations of the best least-squares

 α C(1) is attached to As. The other C atoms are numbered in succession so that $C(4)$ is *para* to $C(1)$. ^b The errors are derived from estimated errors in the group parameters and are meant to be used in error analyses on inter-ring distances. As always, intra-ring distances are fixed by the group refinement procedure.

RNION 1

ANION 2

Figure 1.-The $Co(tdt)_2$ ⁻ anions showing their bond distances and ellipsoids of thermal motion. The ellipsoids have been scaled to 30% probability distributions.

planes through the ligands 21 and the distances of the atoms from their respective planes are given. Only three possibly significant deviations from planarity are noted in the anions.

One of the most interesting and surprising features of the anion geometry is the complete ordering of the tdt methyl groups which necessarily yields a *trans* configuration for the complex. It is possible that the *trans* isomer preferentially crystallizes out from a solution containing equal amounts of both geometric configurations.

The $(C_6H_5)_3(CH_3)As^+$ cation had its expected tetra-**(21' W. C. Hamilton,** *Ada Cvysl.,* **14,** 185 (1961).

hedral geometry about the arsenic. The independent cation distances and angles are listed in Table V. It appears that, because of the asymmetric nature of the $(C_6H_5)_3(CH_3)As⁺$ ion, enough volume remains unoccupied in the unit cell for an ethanol solvent molecule to be trapped during crystallization. The ethanol atoms are represented in Table I as $X(1)$, $X(2)$, and $X(3)$ with $X(1)$ probably being the oxygen. As was previously mentioned, the high-temperature factors for these atoms were taken as an indication that the ethanol molecule sites are not fully occupied. An examination of the final distances calculated for the complex reveals that the ethanol molecules can occupy only half of the solvent molecule sites and that they are probably disordered about the $(0, \frac{1}{2}, 0)$ center of symmetry. Thus, the molecular formula of the complex is $[(C_6H_5)_8(CH_3)As] [Co(tdt)_2] \cdot 0.5C_2H_5OH$.²⁰ An additional cycle of refinement was carried out in which $X(1)$, $X(2)$, and $X(3)$ were assigned fixed multiplicity factors of 0.5. $X(1)$ was also assumed to be the ethanol oxygen. In this final calculation, no anion or cation parameter changed by more than one standard deviation and the conventional and weighted *R* factors remained at their previous values of 0.083 and 0.077. The final parameters obtained for the ethanol molecule from this refinement are given in footnote *b* of Table I and the dimensions of the ethanol molecule are given in Table V. However, because of the observed disorder no true significance should be attached to these dimensions.

The crystal structure of the complex consists of the packing of essentially planar anions, bulky cations, and trapped solvent molecules. The anions are very well separated with the nearest Co-Co distance being the lattice repeat value of 10.21 **8.** The calculated dihedral angle between the two $Co(tdt)₂-$ anions in the unit cell is 60.5° . A stereoview of the packing of the complex is presented in Figure 2. All intermolecular contacts in the crystal appear normal with the closest approaches to the two Co's being made by phenyl hydrogens at distances of 3.34 and 3.28 **8,** respectively. The nearest approach of the ethanol molecule to the planar anions occurs between $X(1)$ and $C(1)$ at approximately 3.5 **8.** An examination of the thermal parameters in Table I also indicates that in the crystal, anion 2 is vibrating more freely than anion 1.

Discussion

The $Co(tdt)₂$ anions in $[(C_6H_5)_3(CH_3)As][Co (tdt)_2$ \cdot 0.5C₂H₅OH are completely separated whereas in the $[(C_6H_5)_3(CH_3)P][Ni(mnt)_2]^7$ and $[(n-C_4H_9)_4N]$ - $[Cu(mnt)₂]$ ⁵ complexes the monoanionic planar units are separated by less that 4.5 Å .²² On the basis of the reasonably close proximity of the anions in the *[(G-* $H₅$ ₃(CH₃)P][Ni(mnt)₂] structure, the low magnetic moment determined for that complex in the solid $(\mu_{eff} \sim 1.0$ BM) was explained by a spin-quenching mechanism in which magnetic interactions were pro-

⁽²²⁾ In the $[(C_6H_8)_8(CH_8)P][Ni(mnt)_2]$ crystal, the observed $Ni \cdot \cdot Ni$ distances are 4.40 and 4.45 Å. The observed $Cu \cdot \cdot Cu$ distances in $[(n-$ CiRp)aNI1Cu(mnt)r] are 4.03 and **4.43 A.**

TABLE I11

OBSERVED AND CALCULATED STRUCTURE FACTOR AMPLITUDES (IN ELECTRONS \times 10) FOR $[(C_6H_5)_3(CH_3)AS][C_0(tdt)_2] \cdot 0.5C_2H_5OH$

TABLE IV

ROOT-MEAN-SQUARE AMPLITUDES OF VIBRATION (IN **A)**

Atom	Minimum	Intermediate	Maximum
Co(1)	0.202(6)	0,228(6)	0.273(6)
Co(2)	0.193(6)	0.253(5)	0.331(7)
As	0.180(3)	0.226(3)	0.242(3)
S(1)	0.209(7)	0.255(6)	0.274(9)
S(2)	0.209(7)	0.249(6)	0.277(9)
S(3)	0.242(7)	0.294(7)	0.346(8)
S(4)	0.240(7)	0.257(7)	0.381(10)

posed between the closely spaced planar units.12 On the other hand, magnetic susceptibility measurements on $Co(tdt)₂$ show the complex to be a full-spin triplet in the solid $(\mu_{eff} = 3.18 \text{ BM})$ with no indication of any pairwise magnetic interactions. **2b** The separation of the $Co(tdt)₂$ anions in the crystal is entirely consistent with these observations.

The assignment of the unusual high-spin groundstate configuration of $Co(tdt)_2$ ⁻ represents an important step in elucidating the electronic structures of the ML_2^{n-} complexes containing the unsaturated bidentate sulfur ligand systems. A molecular orbital calculation of the SCCC type has recently been performed by Hawkes on the complex $Ni(bdt)₂^-$ (bdt = benzene-1,2-dithiolate) assuming D_{2h} symmetry and using the structural parameters from $Co(tdt)₂$ ^{-.23} This calculation, which assumes no hybridization of the sulfur 3s and 3p orbitals, yields an energy level ordering which differs significantly from that obtained in an earlier calculation by Shupack, *et al.,* on the related $Ni(mnt)_2$ ⁻ system.²⁴ In the earlier calculation, sulfur hybrid orbitals were employed for the in-plane σ and π bonding. The energy level ordering proposed by Hawkes²³ gives good agreement with *both* optical and esr data for the $\text{Nil}_2{}^{n-}$ $(n = 0, 1, 2)$ complexes whereas the earlier calculation exhibits inconsistencies with the esr data for the monoanionic systems. **²⁵**

The ordering of the energy levels of interest ob-The ordering of the energy levels of interest obtained by Hawkes is $4a_g \sim 5a_g (x^2 - y^2, z^2) < L(\pi) <$ tained by Hawkes is $4a_g \sim 5a_g (x^2 - y^2, z^2) < L(\pi) < 3b_{3g} (yz) < 4b_{2g}(xy) < 3b_{1g} (xy)$ as opposed to the ordering of $3a_g(z^2)$ < $3b_{3g}(yz)$ < L (π) < $4b_{2g}(xz)$ < $4a_{\mathbf{g}}(x^2 - y^2) < 3b_{\mathbf{1g}}(xy)$ put forth by Shupack, *et al.* Because of the consistency of the former with experimental data and the greater inherent accuracy of the more recent calculation in which no hybridization was assumed, we use the ordering proposed by Hawkes in

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⁽²³⁾ M. J. Hawkes, Ph.D. Thesis, Columbia University, 1967.

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^aBecause of the observed disorder, no true significance should be attached to these values. Estimated standard deviations appear overly op timistic for these dimensions.

 $X(2)-X(3)$ 1.18 (5)

TABLE VI BEST WEIGHTED LEAST-SQUARES PLANES THROUGH THE tdt LIGANDS AND DISTANCES OF THE ATOMS FROM

THEIR RESPECTIVE PLANES [®]					
Atoni	Distance, Å	Atom	Distance, Å		
		Anion 1: $13.852x - 9.416y + 6.728z$	$=-0.100$		
(Triclinic Coordinates)					
S(1)	0.008(5)	C(4)	0.00(2)		
S(2)	$-0.002(5)$	C(5)	$-0.02(2)$		
C(1)	$-0.04(2)$	C(6)	$-0.08(2)$		
C(2)	$-0.02(2)$	M(1)	0.10(2)		
C(3)	$-0.01(2)$				
		Anion 2: $-0.952x + 2.669y - 9.826z = -5.391$			
(Triclinic Coordinates)					
S(3)	$-0.005(5)$	C(10)	0.02(2)		
S(4)	0.004(4)	C(11)	0.03(2)		
C(7)	$-0.03(2)$	C(12)	0.06(2)		
C(8)	$-0.05(2)$	M(2)	0.01(3)		
C(9)	$-0.04(2)$				

^a Weights were based on the variance-covariance matrix as obtained from the final cycle of least-squares refinement.

Figure 2.---These drawings present a view of the packing of $[(C_6H_5)_3(CH_3)As]$ $[Co(tdt)_2] \cdot 0.5C_2H_5OH$ in the crystal. The ethanol molecules are shown in both possible orientations. The drawings constitute a stereopair and may be viewed with a hand viewer.

suggesting a ground-state configuration for the *Co-* $(tdt)_2$ ⁻ complex. The separation of the 3b_{3*u*} (yz) and 4bg *(xz)* orbitals is apparently too small to yield a low-spin $(3b_{3g})^2$ = 1A_g ground-state configuration, and the $3b_{3g}$ (yz) and $4b_{2g}$ (xz) levels are calculated to lie very close to one another. Thus, the predicted ground state is the spin-triplet configuration $(3b_{3g})^1$ $(4b_{2g})^1$ = ${}^3B_{1g}$ and the two unpaired electrons are located in out-of-plane *x* orbitals which are delocalized over the entire complex. A full discussion of the electronic structures of the ML_2^{n-} complexes including details of the calculation, population analyses, and experimental data will be presented in a forthcoming paper.²⁶

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(26) hl. J. **Bakcr-Hawkes** and 13. **13.** Gray, Lo be published.